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On the Hankel Integral Transform with Wright's Generalized Hypergeometric Function and Applications

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Abstract: The aim of this paper is to evaluate Hankel transform of Wright's generalized hypergeometric function defined by Dotsenko [1, 2]. The author has given two applications of Hankel transform of Wright's generalized hypergeometric function by connecting this, first with the Weyl integral and second is with Riemann-Liouville type of fractional derivative. The results obtained are basic in nature and are likely to find useful applications.

(Mathematics Subject Classification: 33C20, 33E20)

Key words: Hankel transform, Wright's Generalized Hypergeometric Function, Weyl Integral, Riemann-Liouville type fractional derivative.

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I. INTRODUCTION

Generalized Wright's function ${}_{2}R_{1}(a,b;c,w;\mu;z)$ defined by Dotsenko [1, 2] hs been denoted as

$${}_{2}R_{1}(a,b;c,w;\mu;z) = \frac{\Gamma(c)}{\Gamma(a)\Gamma(b)} \sum_{k=0}^{\infty} \frac{\Gamma(a+k)\Gamma\left(b+k\frac{w}{\mu}\right)}{\Gamma\left(c+k\frac{w}{\mu}\right)} \frac{z^{k}}{k!}$$

$$= \frac{\Gamma(c)}{\Gamma(a)\Gamma(b)} {}_{2}\psi_{1} \left[z\Big|_{\left(c,\frac{w}{\mu}\right)}^{(a,1),\left(b,\frac{w}{\mu}\right)}\right]$$

$$(1.1)$$

Provided $\operatorname{Re}(c) > 0, \operatorname{Re}(b) > 0, \operatorname{Re}\left(\frac{w}{k}\right) > 0$.

Virchenko et. al. [6] defined the Wright type hypergeometric function by taking $\frac{w}{k} = \tau > 0$ in (1.1) as

$${}_{2}R_{1}^{\tau}(z) = {}_{2}R_{1}(a,b;c,w;\mu;z) = \frac{\Gamma(c)}{\Gamma(b)} \sum_{k=0}^{\infty} \frac{(a)_{k} \Gamma\left(b+k\frac{w}{\mu}\right)}{\Gamma\left(c+k\frac{w}{\mu}\right)} \frac{z^{k}}{k!}, \tau > 0, |z| < 1 \quad (1.2)$$

If $\tau = 1$, then (1.2) reduces to a Gauss's hypergeometric function.

The Hankel transform of a function f(x), denoted by g(p, v) or in short by simply g(p) is defined as

$$g(p;v) = \int_{0}^{\infty} (px)^{\frac{1}{2}} J_{v}(px) f(x) dx; p > 0$$
 (1.3)

Where p as a complex parameter.

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II. HANKEL TRANSFORM OF WRIGHT'S GENERALIZED WRIGHT'S FUNCTION

If $a,b,c,p,A,v \in C$; Re(a) > 0, Re(b) > 0, Re(c) > 0, $Re(\delta) > 0$, $Re(\rho) > 0$, $Re(\nu) > 0$, $Re(\lambda) > 0$ and

$$\begin{split} &\frac{w}{k} \in N \text{ , then } \\ &\int_{0}^{\infty} x^{\rho-1} J_{\nu}(Ax) {}_{2}R_{1}(a,b;c,w;\mu;px^{\delta}) dx \\ &= \frac{\Gamma(c)2^{\rho-1}}{\Gamma(a)\Gamma(b)} {}_{3} \psi_{2} \left[\left(\frac{2^{\delta} p}{A^{\delta}} \right) \Big|_{\left(c,\frac{w}{\mu}\right)}^{\left(a,1\right)} \left(\frac{b,\frac{w}{\mu}}{2},\frac{(\rho+\nu,-\delta)}{2} \right) \right] \quad (2.1) \end{split}$$

$$&\text{Proof: } \int_{0}^{\infty} x^{\rho-1} J_{\nu}(Ax) {}_{2}R_{1}(a,b;c,w;\mu;px^{\delta}) dx \\ &= \int_{0}^{\infty} x^{\rho-1} J_{\nu}(Ax) \frac{\Gamma(c)}{\Gamma(a)\Gamma(b)} \sum_{k=0}^{\infty} \frac{\Gamma(a+k)\Gamma\left(b+\frac{w}{\mu}k\right)}{\Gamma\left(c+\frac{w}{\mu}k\right)} \frac{p^{k}x^{\delta k}}{k!} dx \\ &= \int_{0}^{\infty} x^{\rho+\delta k-1} J_{\nu}(Ax) \frac{\Gamma(c)}{\Gamma(a)\Gamma(b)} \sum_{k=0}^{\infty} \frac{\Gamma(a+k)\Gamma\left(b+\frac{w}{\mu}k\right)}{\Gamma\left(c+\frac{w}{\mu}k\right)} \frac{p^{k}}{k!} dx \end{split}$$

Changing the order of integration and summation therein (which is permissible under the conditions mentioned with (2.1), we find that

$$\begin{split} &= \frac{\Gamma(c)}{\Gamma(a)\Gamma(b)} \sum_{k=0}^{\infty} \frac{\Gamma(a+k)\Gamma\left(b+\frac{w}{\mu}k\right)}{\Gamma\left(c+\frac{w}{\mu}k\right)} \frac{p^{k}}{k!} \int_{0}^{\infty} x^{\rho+\delta k-1} J_{\nu}(Ax) dx \\ &= \frac{\Gamma(c)}{\Gamma(a)\Gamma(b)} \sum_{k=0}^{\infty} \frac{\Gamma(a+k)\Gamma\left(b+\frac{w}{\mu}k\right)\Gamma\left(\frac{\rho+\delta k+\nu}{2}\right)}{\Gamma\left(c+\frac{w}{\mu}k\right)\Gamma\left(1-\frac{-\nu+\rho+\delta k}{2}\right)} \frac{2^{2+\delta k-1}A^{-\rho-\delta k}p^{k}}{k!} \\ &= \frac{\Gamma(c)2^{\rho-1}}{\Gamma(a)\Gamma(b)A^{\rho}} \sum_{k=0}^{\infty} \frac{\Gamma(a+k)\Gamma\left(b+\frac{w}{\mu}k\right)\Gamma\left(\frac{\rho+\delta k+\nu}{2}\right)}{\Gamma\left(c+\frac{w}{\mu}k\right)\Gamma\left(1-\frac{-\nu+\rho+\delta k}{2}\right)} \frac{1}{k!} \left(\frac{2^{\delta}p}{A^{\delta}}\right)^{k} \\ &= \frac{\Gamma(c)2^{\rho-1}}{\Gamma(a)\Gamma(b)} {}_{3}\Psi_{2} \left[\left(\frac{2^{\delta}p}{A^{\delta}}\right)^{\left(a,1\right),\left(b,\frac{w}{\mu}\right),\left(\frac{\rho+\nu-\delta}{2},\frac{-\delta}{2}\right)}}{\Gamma\left(c,\frac{w}{\mu},\frac{\rho+\nu-\delta}{2},\frac{-\delta}{2}\right)}\right] \end{split}$$

III. APPLICATIONS

The Weyl integral ([3], p.91) of f(x) of order α , denoted by W_{∞}^{α} , is defined by

$$\left({}_{x}W_{\infty}^{\alpha}f \right)(x) = \left({}_{x}I_{\infty}^{\alpha}f \right)(x) = \left(I_{-}^{\alpha}f \right)(x)$$

$$= \frac{1}{\Gamma(\alpha)} \int_{x}^{\infty} (x-t)^{\alpha-1} f(t)dt, -\infty, x < \infty$$
(3.1)

Where $\alpha \in C$, Re(α) > 0.

The Weyl Integral of Hankel Transform in Association with Wright's Generalized Hypergeometric Function

The main integral (2.1) can be rewritten as the following Weyl integral formula:

$$\left({}_{0}W_{\infty}^{\rho}J_{\nu}(At){}_{2}R_{1}(a,b;c,w;\mu;pt^{\delta})(t)\right)$$

$$= \frac{1}{\Gamma(\rho)} \int_{0}^{\infty} t^{\rho-1} J_{\nu}(At) {}_{2}R_{1}(a,b;c,w;\mu;pt^{\delta})dt$$

$$= \frac{\Gamma(c)2^{\rho-1}}{\Gamma(\rho)\Gamma(a)\Gamma(b)A^{\rho}} {}_{3}\psi_{2} \left[\left(\frac{2^{\delta}p}{A^{\delta}} \right) \middle|_{\left(c,\frac{w}{\mu}\right)}^{(a,1)} \left(\frac{\rho+\nu-\delta}{2},\frac{-\delta}{2}\right) \right] (3.2)$$

Provided

 $a, b, c, p, A, v \in C$; Re(a) > 0, Re(b) > 0, Re(c) > 0, Re(δ) > 0, Re(ρ) > 0, Re(ν) > 0, Re(λ) > 0 and

$$\frac{w}{k} \in N$$
,

Fractional Derivatives

Following Miller ([4],p.82), let $g \in A$ (Where A is a class of good functions). Then

$${}_{z}D_{\infty}^{q}g(z) = \frac{(-1)^{q}}{\Gamma(-q)} \int_{z}^{\infty} (u-z)^{-q-1} g(u) du, \ for \ q < 0$$
 (3.3)

For $q \ge 0$

$${}_{z}D_{\infty}^{q}g(z) = \frac{d^{r}}{dz^{r}} \left({}_{z}D_{\infty}^{q-r}g(z)\right)$$
(3.4)

r being a positive integer such that r > q.

Fractional Derivatives of Hankel Transform in Association with Wright's Generalized Hypergeometric Function

The main integral (2.1) can be rewritten as the following Riemann-Liouville fractional derivative formula:

$$\frac{1}{2} \int_{-\infty}^{-\rho} \left(J_{\nu}(Ax) {}_{2}R_{1}(a,b;c,w;\mu;px^{\delta}) \right) dx$$

$$= \frac{(-1)^{\rho}}{\Gamma(\rho)} \int_{0}^{\infty} x^{\rho-1} J_{\nu}(Ax) {}_{2}R_{1}(a,b;c,w;\mu;px^{\delta}) dx$$

$$= \frac{(-1)^{\rho} \Gamma(c) 2^{\rho-1}}{\Gamma(\rho)\Gamma(a)\Gamma(b)A^{\rho}} {}_{3}\Psi_{2} \left[\left(\frac{2^{\delta} p}{A^{\delta}} \right) \left| \frac{(a,1) \left(b, \frac{w}{\mu} \right) \left(\frac{\rho+\nu}{2}, \frac{-\delta}{2} \right)}{\left(c, \frac{w}{\mu} \right) \left(1 - \frac{\rho-\nu}{2}, \frac{-\delta}{2} \right)} \right| ; for \rho > 0. \quad (3.5)$$

It is being assumed that the conditions given in (2.1) and (3.3) are satisfied.

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