

Ultraspherical Solutions for Neutral Functional Differential Equations

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ABSTRACT

This paper is concerned with the numerical solution of neutral functional differential equations (NFDEs). Based on the ultraspherical v-stage continuous implicit Runge-Kutta method is proposed. The description and outlines algorithm of the method are introduced. Numerical results are included to confirm the efficiency and accuracy of the method.

Keywords: Functional differential equations, Equations of neutral type, Implicit delay equations.

I. INTRODUCTION

We are interested in the numerical solution of initial-value problem for neutral functional differential equations (NFDEs), which take the form:

$$y'(t) = f(t, y(.), y'(.)) t_0 < t \le t_N, ,$$

$$y(t) = g(t) \alpha \le t \le t_0 (1)$$

where $g(t) \in C^{1}[\alpha, t_{0}]$ and the function f(t, y(.), y'(.)) satisfies the following conditions:-

 $H_{1}: \text{ For any } y \in C^{1}[t_{0}, t_{N}] \text{ the mapping } t \rightarrow f(t, y(.), y'(.)) \text{ is continuous on}$ $[t_{0}, t_{N}].$

 H_2 : There exist constants $L_1 \ge 0$, $0 \le L_2 \le 1$ such that

$$\left\| f(t, y_{1}(.), z_{1}(.)) - f(t, y_{2}(.), z_{2}(.)) \right\| \leq \Lambda$$
$$\Lambda = L_{1} \left\| y_{1} - y_{2} \right\|_{C^{1}[\alpha, t_{N}]} + L_{2} \left\| z_{1} - z_{2} \right\|_{C^{0}[\alpha, t_{N}]}$$

for any $t \in [t_0, t_N]$, y_1 , $y_2 \in C^1[\alpha, t_N]$ and z_1 , $z_2 \in C^0[\alpha, t_N]$.

Under the conditions H_1 and H_2 the problem (1) has a unique solution y(x) [1]. The equations of type (1) have applications in many fields such as control theory, oscillation theory, electrodynamics, biomathematics, and medical science. Numerical methods for the problem (1) were discussed extensively by many authors; see [2-26].

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This paper is concerned with the numerical solution of neutral functional differential equations (NFDEs). Based on the ultraspherical ν -stage continuous implicit Runge-Kutta method is proposed. In section 2, we will adapt a finite

Ultraspherical expansion to approximate $\int_{I_i} f(s) ds$ and f(.) on the interval I_i , i = 0 (1) N - 1. Also,

an easily implemented numerical method for NFDEs will be derived. Finally, in section 3 we present some numerical examples; which show that the presented method provides a noticeable improvement in the efficiency over some previously suggested methods.

II. THE NUMERICAL METHOD

2.1 The Description of the method

Let $\Delta := (t_0 < t_1 < \cdots < t_N)$ define a partition for $[t_0, t_N]$, with the step size $h_i = t_{i+1} - t_i$.

Each subinterval $I_i = [t_i, t_{i+1}]$ is divided by the Chebyshev collocation points:

$$t_{ij} = t_i + h_i \xi_j, \quad \xi_j = \frac{1}{2} (1 - \cos \frac{j\pi}{v}), \quad j = 0 (1)v$$
 (2)

This method is based on a finite Ultraspherical expansion in each subinterval I_i , i = 0 (1) N - 1. Consider the approximation f(t) of f(t) for $t = t_i + \theta h_i$, $0 \le \theta \le 1$, i = 0 (1) N - 1 as follows:

$$\tilde{f}(t) = \sum_{r=0}^{\nu} a_r C_r^{[\alpha]} \left(\frac{2(t-t_i)}{h_i} - 1 \right), t \in [t_i, t_{i+1}]$$
(3)

where

$$a_{r} = \frac{\pi}{v \lambda_{r}^{[\alpha]}} \sum_{l=0}^{v} \int_{0}^{\pi} f(t_{il}) \left(\sin\left(\frac{l\pi}{v}\right) \right)^{2\alpha} C_{r}^{[\alpha]} \left(-\cos\left(\frac{l\pi}{v}\right) \right)$$
$$t_{il} = t_{i} + \frac{h_{i}}{2} \left(1 - \cos\left(\frac{l\pi}{v}\right)\right)$$
(4)

Here $C_r^{[\alpha]}(x)$ is the r-th Ultraspherical polynomial. As especial cases, at $\alpha = 0$ give Chebyshev Polynomials of the first kind $C_r^{[0]}(x) = T_r(x)$, at $\alpha = \frac{1}{2}$ give Legendre Polynomials $C_r^{\left[\frac{1}{2}\right]}(x) = P_r(x)$, at $\alpha = 1$ give Chebyshev Polynomials of the second kind $C_r^{\left[1\right]}(x) = \frac{1}{1+r}R_r(x)$. A summation symbol with double prims denotes a sum with the first and last terms halved.

Now, we can easily show that the following relations are true:

$$I_{r}(\theta) = \int_{-1}^{2\theta - 1} C_{r}^{[\alpha]}(s) \, ds = \begin{cases} C_{1}^{[\alpha]}(2\theta - 1) + 1 & \text{if } r = 0\\ \frac{(2\alpha + 1)}{4(\alpha + 1)} [C_{2}^{[\alpha]}(2\theta - 1) - 1] & \text{if } r = 1 \end{cases} \\ \frac{1}{2(\alpha + r)} (\frac{2\alpha + r}{r + 1} \Phi_{1} - \frac{r}{2\alpha + r - 1} \Phi_{2}) & \text{if } r \ge 2 \end{cases}$$

where
$$\Phi_1 = C_{r+1}^{[\alpha]} (2\theta - 1) + (-1)^r$$
, $\Phi_2 = C_{r-1}^{[\alpha]} (2\theta - 1) + (-1)^r$

Using (3), (4) and (5) the indefinite integral $\int_{t_i}^{t_i + \theta + h_i} \int_{t_i}^{\infty} f(s) ds$ takes the form

$$\int_{t_{i}}^{t_{i}+\theta h_{i}} \tilde{f}(s) d s = \frac{h_{i}}{2} \sum_{r=0}^{\nu} b_{l}^{[\alpha]}(\theta) \tilde{f}(t_{il})$$
(6)

where

$$b_{l}^{\left[\alpha\right]}(\theta) = \frac{\pi \left[\sin\left(\frac{l\pi}{\nu}\right)\right]^{2\alpha}}{2\left(1+\delta_{l0}+\delta_{l\nu}\right)} \sum_{l=0}^{\nu} \frac{I_{r}(\theta) C_{r}^{\left[\alpha\right]}\left(-\cos\left(\frac{l\pi}{\nu}\right)\right)}{\lambda_{r}^{\left[\alpha\right]}}, l = 0 (1) \nu$$
(7)

and δ_{lv} is the Kronecker delta.

From the relations (2), (6) and (7), we obtained the following results:

$$\begin{bmatrix} t_{ij} & \tilde{f}(s) & ds \end{bmatrix} = \frac{h_i}{2} B^{\lceil \alpha \rceil} F_i$$
(8)

$$\begin{bmatrix} t_{ij} & t_{ij} & t_{ij} \\ \int \\ f_{i} & \cdots & f_{i} \\ \vdots & \vdots & t_{i} \end{bmatrix} = \left(\frac{h_{i}}{2}\right)^{k} \quad \left(B^{\left[\alpha\right]}\right)^{k} \quad F_{i} = \frac{\left(\frac{h_{i}}{2}\right)^{k}}{(k-1)!} \quad B^{\left[\alpha\right]} \quad K \quad F_{i}$$
(9)

$$\int_{t_{i}}^{t_{ij}} \tilde{f}(s,t) ds = \sum_{l=0}^{i-1} \frac{h_{l}}{2} \sum_{s=0}^{\nu} b_{\nu s}^{[\alpha]} \tilde{f}(t,t_{ls}) + \frac{h_{i}}{2} \sum_{s=0}^{\nu} b_{js}^{[\alpha]} \tilde{f}(t,t_{is}) , j = 0 (1)\nu$$
(10)
where

where

$$B^{[\alpha]} = [b_{js}^{[\alpha]}]_{j,s=0}^{\nu}, b_{js}^{[\alpha]} = b_{s}^{[\alpha]}(c_{j}), F_{i} = [f(t_{i0}), \cdots, f(t_{i\nu})]^{T}$$

and $K = diag [(c_j - c_0)^{k-1}, \cdots, (c_j - c_v)^{k-1}].$

We notice that the matrix B of El-Gendi's method [27], is the same matrix $B^{\left[\alpha\right]} = \begin{bmatrix} 2 & b_{js}^{\left[0\right]} \end{bmatrix}_{j,s=0}^{\nu}$. On the interval $(t_i, t_{i+1}]$, rewriting (1) in the form:

$$\tilde{y}(t) = \tilde{y}(t_i) \int_{t_i}^t \tilde{z}(s) ds, \ \tilde{z}(t) = f(t, \tilde{y}(.), \tilde{z}(.))$$
 (11)

Suppose the approximations of y(t) and z(t) are given for $t \le t_i$. On $(t_i, t_{i+1}]$, we define the ν -stage method $[\nu - UM]$, as follows

$$z(t_{il}) = f(t_{il}, y(.), z(.)) , l = 1(1)v$$
(12)

where

$$\tilde{y}(t) = \begin{cases} g(t) & \text{if } t \leq t_0 \\ \tilde{y}(t_{j-1}, v) + \frac{h_j}{2} b_0^{[\alpha]}(\theta) & \tilde{z}(t_{j-1}, v) + \frac{h_j}{2} \sum_{l=1}^{v} b_l^{[\alpha]}(\theta) & \tilde{z}(t_{jl}) & \text{if } t = t_j + \theta h_j \in I_j \end{cases}$$
(13)

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$$\tilde{z}(t) = \begin{cases} g'(t) & \text{if } t \leq t_0 \\ \sum_{r=0}^{\nu} a_r C_r^{[\alpha]} (\frac{2(t-t_j)}{h_j} - 1) & \text{if } t = t_j + \theta h_j \in I_j \end{cases}$$
(14)

$$a_{r} = \frac{\pi}{v \lambda_{r}^{\left[\alpha\right]}} \sum_{l=0}^{v} {\overset{''}{z}} \left(t_{jl}\right) \left(\sin\left(\frac{l\pi}{v}\right)\right)^{2\alpha} C_{r}^{\left[\alpha\right]} \left(-\cos\left(\frac{l\pi}{v}\right)\right)$$
(15)

2. 2 The Algorithm of V -stage method

The described method represents a generalization of the methods given by Jackiewicz [7-8], for $\nu = 1,2,3$ and $\alpha = 0$. The algorithm of the ν -stage method is given below:-

STEP 1: Input $N, h_j, I_j, j = 0$ (1) N - 1.

STEP 2: Put i = 0

STEP 3: Compute $y(t_{ij})$ and $z(t_{ij})$ on I_j , j = 1 (1) N - 1, by solving the system of v -equation (12). **STEP 4:** Store the computed values $\tilde{y}(t_{ij})$ and $\tilde{z}(t_{ij})$ on I_j , j = 1 (1) N - 1.

STEP 5: If i = N - 1 go to step 6, otherwise set i = i + 1 and go to step 3.

STEP 6: Output the results $y(t_{ij})$ and $z(t_{ij})$, i = 0 (1) N - 1.

III. NUMERICAL EXAMPLES

In this section, we present the result of some computational experiments by applying our v - UM method.

Example 1: (Jackiewicz [10])

$$y'(t) = e^{-y(t)} + \sin \left(y'(\beta(t)) \right) - \sin \left(\frac{1}{3 + \beta(t)} \right), \quad t \in (0, 10],$$

$$y(0) = \ln (3) , \quad y'(0) = \frac{1}{3}$$
 (16)

Here $\beta(t) = 0.5 t (1 - \cos(2\pi t))$. The exact solution is $y(t) = \ln(3 + t)$.

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Example 2: (Jackiewicz [10])

$$y'(t) = \frac{\sin(t y'(t)) - \sin(e^{y(t)})}{16} + \frac{1}{t}, \quad t \in (1, 4],$$

$$y(t) = 0, \quad y'(t) = 1$$
(17)

Here the exact solution is $y(t) = \ln(t)$.

		Table 1	Tał	ole 2
	Example	1: $y(10) = 2.5649494$	Example	2: $y(4) = 1.386294361$
h	2 - UM	C_0C_1	2 - UM	C_0C_1
2^{-1}	1.21E-4	-2.5E-4	-1.12E-4	-4.00E-1
2^{-2}	7.62E-6	-6.25E-2	-8.00E-6	-7.00E-2
2^{-3}	4.76E-7	6.25E-2	5.00E-7	-5.00E-3
2^{-4}	2.98E-8	7.42E-2	-3.13E-8	3.13E-3
2^{-5}	1.86E-9	2.93E-3	-1.95E-9	-3.13E-4
2^{-6}	1.16E-10	-7.32E-4	-1.22E-10	-3.13E-4
2^{-7}	1.02E-11	-7.93E-4	-7.63E-12	-8.79E-5
2^{-8}	-7.10E-12	-3.66E-4	-4.77E-13	-2.44E-5
2^{-9}	1.15E-11	-8.39E-5	-2.98E-14	-6.10E-6
2^{-10}	-4.88E-12	-5.72E-6	-4.66E-15	-1.53E-6

Example 3: (Jackiewicz [10])

$$y'(t) = e^{1-2t^{2}} y(t^{2}) \left[y'(\frac{t-1}{t+1}) \right]^{(1+t)}, \quad t \in (0,1],$$

$$y(t) = e^{t}, \quad t \in [-1,0]$$
(18)

Here the exact solution is $y(t) = e^{t}$.

Example 4: (Jackiewicz [10])

$$y'(t) = 2\cos(2t)\left(y(\frac{t}{2})\right)^{2\cos(t)} + \ln\left(\frac{y'(\frac{t}{2})}{2\cos(t)}\right) - \sin(t), \quad t \in (0,1],$$

$$y(0) = 1, \quad y'(0) = 2$$
(19)

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	r	Fable 3	Tab	le 4
	Example	3: $y(1) = 2.7182818$	Example	4: $y(1) = 2.4825777$
h	2 - UM	C_0C_1	2 - UM	C_0C_1
2^{-1}	1.25E-3	2.75E-2	-5.94E-2	4.28E-1
2^{-2}	7.81E-5	-2.50E-3	-4.41E-3	5.44E-2
2^{-3}	4.88E-6	-2.03E-3	-3.08E-4	1.91E-2
2^{-4}	3.05E-7	-6.64E-4	-2.11E-5	5.59E-3
2^{-5}	2.86E-8	-1.86E-4	-1.42E-6	1.33E-3
2^{-6}	1.79E-9	-4.88E-5	-9.30E-8	3.13E-4
2^{-7}	7.45E-11	-1.22E-5	-5.96E-9	7.45E-5
2^{-8}	9.31E-12	-3.20E-6	-3.73E-10	1.82E-5
2^{-9}	7.28E-13	-8.01E-7	-2.33E-11	4.46E-6
2^{-10}	-2.73E-15	-2.00E-7	-1.40E-12	1.12E-6

Here the exact solution is $y(t) = e^{\sin(2t)}$

In tables (1)-(4), we give **E** for **2-UM** method and **E** for the best method of Jackiewicz; $C_0 C_1$ method [10], where **E** denotes the global error at the end point t_N .

[10], where **E** denotes the global error at the end point γ_N

Example 5: (Kappel-Kunish [15] and Jackiewicz [10])

$$y'(t) = y(t) + y(t-1) - 0.25 \quad y'(t-1) , t \in (0,2],$$

$$y(t) = -t , \qquad y'(t) = -1, \quad t \in [-1,0]$$
(20)

Here the exact solution is given by

$$y(t) = t + 0.25 \quad e^{t} - 0.25 \quad , \qquad t \in [0,1]$$
$$y(t) = -t + 0.25 \quad e^{t} + 0.5 \quad + \frac{17 + 3t}{16} \quad e^{(t-1)} \quad , \qquad t \in [1,2]$$

Here, the first derivative has a discontinuity at t = 0 and t = 1.

In example 5, we give **E** for the method **2-UM**, **5-UM**, **8-UM** and $C_0 C_1$ in Table (5), the V - **UM** methods, for large V, make a little improvement in the computed results, the reason is due to there exist discontinuity for the first derivative of Y(t) at t = 0 and t = 1. These results indicate that the V - **UM** method is better than the one-step methods of Jackiewicz [10].

	Table 5				
	Exampl	e 5: $y(2) =$	= 4.2547942	2	
h	2 - UM	5 - UM	8 - UM	C_0C_1	
2^{-1}	1.6	2.03E-1	8.00E-2	6.53	
2^{-2}	8.15E-1	1.00E-1	4.00E-2	4.22	
2^{-3}	4.10E-1	4.95E-2	2.00E-2	2.35	
2^{-4}	2.05E-1	2.48E-2	1.00E-2	1.22	
2^{-5}	1.03E-1	1.23E-2	5.00E-3	6.14E-1	
2^{-6}	5.13E-2	6.16E-3	2.50E-3	3.08E-1	
2^{-7}	2.56E-2	3.08E-3	1.25E-3	1.54E-1	
2^{-8}	1.28E-2	1.54E-3	6.25E-4	7.69E-2	
2^{-9}	6.41E-3	7.70E-4	3.13E-4	3.84E-2	
2^{-10}	3.20E-3	7.70E-4	1.56E-4	1.92E-2	

Example 6: (Castleton-Grimm [4] and Jackiewicz [8])

$$y'(t) = \frac{-4 t y^{2}(t)}{\left(\ln (\cos (2 t))\right)^{2} + 4} + \tan (2 t) + \frac{\tan^{-1}(z(t))}{2}, \quad t \in (0, 0.75],$$

$$y(0) = 0, \qquad y'(0) = 0$$
(21)
where $z(t) = y'\left(\frac{t y^{2}}{1 + y^{2}(t)}\right)$. The theoretical solution is $y(t) = -0.5$ ln (cos(2 t)).

Example 7: (Castleton-Grimm [4] and Jackiewicz [8])

$$y'(t) = (1 + u(t)) \cos (t) + y(t) z(t) - \sin (t + t \sin^{2}(t)), \quad t \in (0,1],$$

$$y(0) = 0, \qquad y'(0) = 1$$
(22)

where $u(t) = y(t y^{2}(t))$, $z(t) = y'(t y^{2}(t))$.

The theoretical solution is $y(t) = \sin(t)$.

Example 8: (Pouzet [16] and Jackiewicz [8])

$$y'(t) = e^{-t^{2}}(1+t^{2}) + t - y(t) + \int_{0}^{t} t^{2} e^{-st} y(s) y'(s) ds, y(0) = 0, t \in (0,2]$$
(23)

The exact solution is y(t) = t.

Table 6 : Example 6 (Castleton-Grimm [4] and Jackiewicz [8])					
t_n	$y(t_n)$	$y_n(2^{-4})$	$y_n(2^{-6})$	$y_n(2^{-8})$	$y_n(2^{-10})$
0.0625	0.003916465	0.003916470	0.003926723	0.003916465	0.003917107
		0.003916465	0.003916505	0.003916465	0.003916468
0.1250	0.015790526	0.015790547	0.015832887	0.015790526	0.015793178
		0.015790526	0.015790691	0.015790526	0.015790536
0.1875	0.036012506	0.036012544	0.036113197	0.036012505	0.036018812
		0.036012506	0.036012900	0.036012506	0.036012531
0.2500	0.065292120	0.065292132	0.065486194	0.065292114	0.065304276
		0.065292120	0.065292880	0.065292120	0.065292168
0.3125	0.104766351	0.104766181	0.105105260	0.104766333	0.104787582
		0.104766350	0.104767677	0.104766351	0.104766433
0.3750	0.156199949	0.156199327	0.156765982	0.156199923	0.156235405
		0.156199949	0.156202162	0.156199949	0.156200087
0.4375	0.222365362	0.222364313	0.223301956	0.222365342	0.222423942
		0.222365362	0.222369022	0.222365362	0.222365591
0.5000	0.307813235	0.307813411	0.309395557	0.307813221	0.307912401
		0.307813235	0.307819432	0.307813235	0.307813622
0.5625	0.420618861	0.420623077	0.423464892	0.420618861	0.420797810
		0.420618860	0.420630041	0.420618861	0.420619559
0.6250	0.577079895	0.577100894	0.582906497	0.577079967	0.577448783
		0.577079896	0.577102963	0.577079895	0.577081337
0.6875	0.818538930	0.818682074	0.834410205	0.818539584	0.819566071
		0.818538933	0.818603259	0.818538930	0.818542950
0.7500	1.324391827	1.328895912	1.430890306	1.324428196	1.332263471
		1.324391983	1.324892390	1.324391828	1.324423149

	Table 7 :Exa	ample 7 (Castle	eton-Grimm [4] and Jackiewi	cz [8])
	$y(t_n = 0.25) = 0.247403959,$ $y(t_n = 0.50) = 0.479425539$				
	$y(t_n = 0.75)$	= 0.681638760), $y(t_n =$	= 1.0) = 0.8414	70985
t_n	$y_n(2^{-2})$	$y_n(2^{-4})$	$y_n(2^{-6})$	$y_n(2^{-8})$	$y_n(2^{-10})$
0.25	0.247404002	0.245934360	0.247403959	0.247306957	0.247403959
	0.247397577	0.247403959	0.247403556	0.247403959	0.247403934
0.50	0.479423742	0.475355054	0.479425535	0.479158061	0.479425539
	0.479408626	0.479425539	0.479424471	0.479425539	0.479425472
0.75	0.681635238	0.672679530	0.681638752	0.681051318	0.681638760
	0.681602967	0.681638760	0.681636491	0.681638760	0.681638618
1.00	0.841459232	0.825030283	0.841470963	0.840379343	0.841470984
	0.841404641	0.841470985	0.841466775	0.841470985	0.841470722

	Table 8 :Example 8 (Pouzet [16] and Jackiewicz [8])				
t_n	$y(t_n)$	$y_n(0.1)$	$y_n(0.05)$	$y_n(0.025)$	$y_n(0.0125)$
0.1	0.100000000	0.100000000	0.100000000	0.100000000	0.100000000
		0.099999999	0.0999999999	0.100000000	0.100000000
0.2	0.200000000	0.200000000	0.200000000	0.200000000	0.200000000
		0.199999855	0.199999972	0.19999993	0.199999998
0.3	0.300000000	0.300000000	0.300000000	0.300000000	0.300000000
		0.299999115	0.299999806	0.299999958	0.299999988
0.4	0.400000000	0.3999999999	0.400000000	0.400000000	0.400000000
		0.399996708	0.399999241	0.399999815	0.399999954
0.5	0.500000000	0.499999998	0.500000000	0.500000000	0.500000000
		0.499990890	0.499997844	0.499999469	0.499999868
0.6	0.600000000	0.599999994	0.600000000	0.600000000	0.600000000
		0.599979283	0.599995019	0.599998768	0.599999693
0.7	0.700000000	0.699999989	0.699999999	0.700000000	0.700000000
		0.699959044	0.699990053	0.699997533	0.699999384
0.8	0.800000000	0.799999981	0.7999999999	0.800000000	0.800000000
		0.799927095	0.799982170	0.799995569	0.799998894
0.9	0.900000000	0.899999953	0.899999998	0.900000000	0.900000000
		0.899880395	0.899970604	0.899992684	0.899998173
1.0	1.000000000	0.999999953	0.999999997	1.000000000	1.000000000
		0.999816214	0.999954664	0.999988706	0.999997179
1.1	1.100000000	1.099999933	1.099999996	1.100000000	1.100000000
		1.099732385	1.099933803	1.099983497	1.099995877
1.2	1.200000000	1.199999907	1.199999994	1.200000000	1.200000000
		1.199627509	1.199907665	1.199976968	1.199994245
1.3	1.300000000	1.299999876	1.299999992	1.300000000	1.300000000
		1.299501081	1.299876121	1.299969086	1.299992275
1.4	1.400000000	1.399999839	1.399999990	1.399999999	1.400000000
		1.399353516	1.399839272	1.399959876	1.399989973
1.5	1.500000000	1.499999794	1.499999987	1.4999999999	1.500000000
		1.499186056	1.499797427	1.499949416	1.499987358
1.6	1.600000000	1.599999740	1.599999984	1.5999999999	1.600000000
		1.599000583	1.599751058	1.599937824	1.599984460
1.7	1.700000000	1.699999675	1.699999980	1.699999999	1.700000000
		1.698799343	1.699700726	1.699925339	1.699981314
1.8	1.800000000	1.799999596	1.799999975	1.799999998	1.800000000
		1.798584660	1.799647014	1.799911809	1.799977956
1.9	1.900000000	1.899999502	1.899999969	1.899999998	1.900000000
		1.898358666	1.899590455	1.899897665	1.899974420
2.0	2.000000000	1.999999389	1.999999962	1.999999998	2.000000000
		1.998123103	1.999531488	1.999882918	1.999970733

In Tables (6)-(8), we give the exact solution $y(t_n)$ in the second column, and $y_n(h)$ which denotes

the computed value $y_h(t_n)$ at every $t = t_n$, the computed value of $y_n(h)$ are represented in two lines, the first line for 2-UM and the second line for the method of Jackiewicz [8]. The results given in Tables (6) and (7) are much better than those obtained by Castleton-Grimm [4] and also, those obtained by Jackiewicz [8]. The results given in Table (8) are much better than those obtained by Jackiewicz [8].

When solving the nonlinear equations, the computations are terminated when two successive

approximations differed by less than $10 h_3$.

IV. CONCLUSIONS

In this paper we construct a method based on the Ultraspherical approximation. This method can be applied to solve different types of NFDEs. The experimental comparison, presented in this paper, shows that this method is more efficient than the previously introduced methods. In addition, the V - UM method can be easily implemented on computer compared with the Lagrange multipliers and their integrals which given by Jackiewicz [10].

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